

Negative-positive monopole transitions in cholesteric liquid crystals

M. V. Kurik and O. D. Lavrentovich

Institute of Physics, Academy of Sciences of the Ukrainian SSR

(Submitted 6 March 1982)

Pis'ma Zh. Eksp. Teor. Fiz. **35**, No. 9, 362-365 (5 May 1982)

The nature of the changes in the monopole structures accompanying a transition out of a levorotary into a dextrorotary cholesteric is clarified experimentally for the first time. The concept of “+” and “-” monopoles in cholesterics is introduced. In the inversion region, textures are discovered with equatorial disclinations and surface and volume point defects. Analogies with textures of $^3\text{He-A}$ in a sphere are noted.

PACS numbers: 61.30.Eb, 61.30.Gd, 61.30.Jf

The possibility of the formation of spherical concentric systems of cholesteric liquid crystal (CLC) layers with a defect line emerging from the center of each such system was established in Refs. 1 and 2 for the case $|q|R \sim 10$ ($|q| = 1/P$, where P is the pitch of the CLC helix, and R is the radius of the drop) as a result of investigations of spherical drops of CLC. In the continuum approximation ($|q|R \gg 1$), as shown theoretically in Ref. 3, such textures represent monopoles in CLS, the analog of the Dirac monopole. In addition, both structures with one $\chi(+2)$ and two $\chi(+1)$ lines were predicted. The experimental investigation described in Refs. 4 and 5, carried out for $|q|R \sim 100$, confirmed the results in Ref. 3.

A monopole in CLC is characterized by the presence of a point singularity in the field of the vector \mathbf{n} normal to the cholesteric layers, out of which emerges the χ line.³ Depending on the right or left-handedness of CLC helix (CLCR and CLCL, respectively), the vector \mathbf{n} turns out to be parallel or antiparallel to the radius vectors of the spheres; in other words, it is possible to distinguish “+” and “-” monopoles. In this connection, it is interesting to investigate the nature of the possible transformations of “-” into “+” monopoles and vice versa, which was the subject of this work. Such a transformation can be realized for a mixture of two CLC with different signs of the helicity (sign of q), since it is well known that in such systems at a certain temperature a transition occurs from negative to positive values of q .⁶ At the inversion point, $q = 0$ and the mixture is a nematic liquid crystal (NLC).

In this work, we used a mixture of cholesteryl chloride and cholesteryl myristate in a proportion 1.75:1.00 by weight with an inversion temperature of 41 °C, above which the mixture is a CLCL ($q < 0$) and below which it is a CLCR ($q > 0$). The substance was dispersed in the form of spherical drops with radius $R \lesssim 50 \mu\text{m}$ in isotropic liquids: pure glycerine in order to create tangential boundary conditions and glycerine with additions (up to 1%) of a lecithin solution in order to create normal boundary conditions.⁷ The specimens were studied with the help of an NU-2E polarization microscope and the temperature was varied at a rate 0.2 °C/min.

The results of the investigations of freshly prepared mixtures for transitions of the isotropic liquid (IL) \rightarrow CLCL \rightarrow NLC \rightarrow CLCR are shown in Figs. 1 and 2. We note that for the opposite transitions, the nature of the textural transformations does not change.

Tangential boundary conditions. In an IL-CLCL transition in drops, concentric spherical layers of CLC with χ lines form. Since the mixture at temperatures $t > 41^\circ\text{C}$ is levorotatory, the textures represent “-” monopoles. In addition, the probability for the formation of a $\chi(+2)$ line in a drop is greater than the probability for the formation of two $\chi(+1)$ lines by a factor of approximately 100. As the temperature decreases, the CLCL helix unwinds (Fig. 1a), while at the same time the positioning of the layers remains the same up to values $|q|R \sim 1$. In addition, only drops with $\chi(+2)$ lines are observed. The t dependence of q , measured by using the procedure in Ref. 2, is almost linear, consistent with the data in Ref. 6. The size of the cores of the $\chi(+2)$ lines increases with an increase in the pitch of the helix. For $|q|R \lesssim 1$, at the location where the χ lines reach the surface of the drop, the core degenerated into two singular points (Fig. 1b), which separated toward the poles of the drop (Fig. 1c); as a result, a so-called bipolar texture² with two diametrically opposite point surface singularities of strength +1 each were established. With further decrease in t and inversion of the sign of q , the singulari-

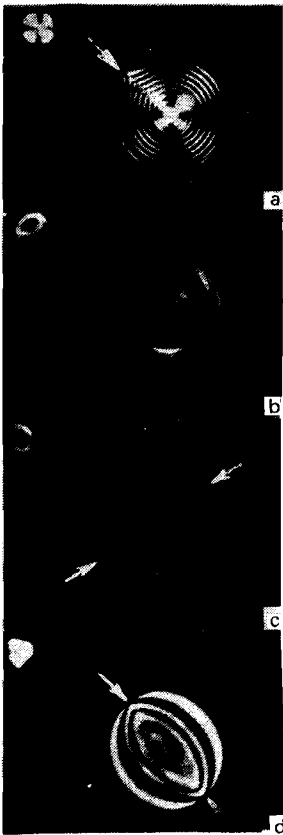


FIG. 1. CLC drop with $R = 35 \mu\text{m}$ in glycerine. The nichols are crossed. a) Monopole structure with $\chi(+2)$ line, $P = 6.5 \mu\text{m}$; b, c, d) show the degeneration of the $\chi(+2)$ line into two singular points and their divergence toward the poles of the drop ($P \approx 40, 60 > 100 \mu\text{m}$, respectively). The defects are indicated with arrows.

ties converge and, at the point where they are localized, $\chi(+2)$ disclinations of "+" monopoles with a dextrorotary helix emerge.

Thus, for CLC drops with tangential boundary conditions, transitions of "+" into "-" monopoles (and vice versa) are caused by mutual transformations of $\chi(+2)$ disclinations and surface point singularities with intensity +1. We note that the convergence of two points of a bipolar texture is reminiscent of a transition to the boojum type texture⁸ with a single surface singularity with intensity +2, but only for the case $|q|R < 1$, since for $|q|R \gg 1$, as pointed out in Ref. 5, the formation of a boojum is related to considerable energy losses due to the existence of the condition for equidistant separation of CLC layers.³

Normal boundary conditions. For high values $|q|R > 10$, the textures in the drops represent "-" monopoles; i.e., in this case, the conditions on the surface do not change the packing of the layers. However, they are manifested for $|q|R < 5$: the appearance of an additional defect, an equatorial disclination, is observed (Figs. 2a and 2b). Such a defect was predicted theoretically⁹ and observed experimentally² for drops of a nematic liquid crystal in a strong magnetic field. The results of this study indicate that the presence of an external magnetic field is not a necessary condition for the formation of a sur-

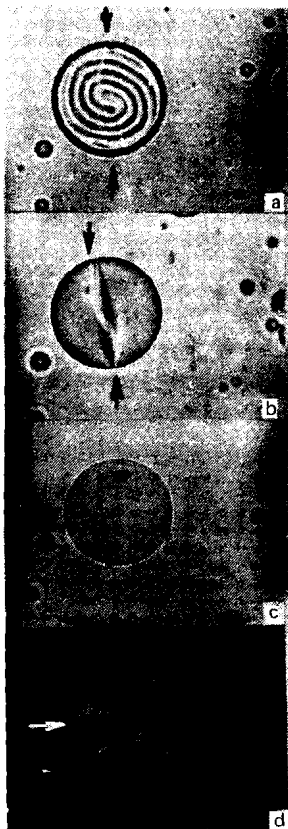


FIG. 2. The CLC drop with $R = 38 \mu\text{m}$ in glycerine but with the additions of lecithin. The nichols are missing. a) Equatorial disclination, $P = 16 \mu\text{m}$; b) equatorial disclination, $P = 80 \mu\text{m}$; c) point surface singularity, $P \geq 80 \mu\text{m}$; d) same, in crossed nichols. The defects are indicated with arrows.

face disclination. In the region $|q|R < 1$, a disclination can relax into a point surface defect (Fig. 2c). Investigation of this texture in crossed nicols (Fig. 2d) confirms that the configuration of the director \mathbf{d} for it is analogous to the distribution of the vector \mathbf{l} of the orbital angular momentum for superfluid $^3\text{He-A}$ in the boojum texture.^{8,10,11} The contraction of the disclination into a point defect is analogous to the relaxation of the texture of $^3\text{He-A}$ in a sphere with equatorial semidisgyration to the texture of a boojum^{10,11} and results from the energy conditions: the energy of linear defects, in contrast to point defects, contains logarithmically divergent terms $\sim R \ln R$.

Textures with equatorial disclinations and surface point defects were not the only possible ones: Their transitions into radial structures (studied in detail in Ref. 2) and the reverse transitions out of the radial textures into textures with surface point defects were observed.

A further decrease in temperature results in the formation of "+" monopoles. In addition, if these monopoles are preceded by textures with a surface point singularity, $\chi(+2)$ lines will emerge at the location of this point.

Thus, the transition "+" to "-" monopoles in CLC (and vice versa) under normal boundary conditions is characterized by the formation of additional surface defects, equatorial disclinations, and a point singularity analogous to the singularity of the \mathbf{I} vector in the boojum structure for $^3\text{He-A}$, as well as the formation of a point defect in the drop.

1. C. Robinson, *Trans. Far. Soc.* **52**, 517 (1956).
2. S. Candau, P. Le Roy, and F. Debeauvais, *Mol. Cr. Liq. Cr.* **23**, 283 (1973).
3. G. E. Volovik, *Pis'ma Zh. Eksp. Teor. Fiz.* **29**, 357, 522 (1979) [*JETP Lett.* **29**, 322 (1979)].
4. M. V. Kurik and O. D. Lavrentovich, Abstracts of Reports at the IV-th Intern. Conf. of Socialist Countries on Liquid Crystals (Tbilisi, 5-8 October, 1981), Pt. 1, p. 173.
5. M. V. Kurik and O. D. Lavrentovich, *Mol. Cr. Liq. Cr. Lett.* (in press).
6. E. Sackmann, S. Meiboom, L. C. Snyder, A. E. Meixner, and R. E. Dietz, *J. Am. Chem. Soc.* **90**, 3567 (1968).
7. A. J. Derzhanski, S. B. Naidenova, and B. Tavena, *Compyt. Rend. Acad. Bulg. Sc.* **31**, 517 (1978).
8. D. L. Stein, R. D. Pisarski, and P. W. Anderson, *Phys. Rev. Lett.* **40**, 1269 (1978).
9. E. Dubois-Violette and P. Parodi, *J. Phys.* **30**, Suppl. 57 (1969).
10. N. D. Mermin, *Quantum Fluids and Solids*, Ed. by S. B. Trickey, E. D. Adams, and J. F. Dufty, Plenum Press, N. Y., 1977, p. 3.
11. P. W. Anderson and R. G. Palmer, *ibid.* p. 23.

Translated by M. E. Alferieff

Edited by S. J. Amoretty